

No [CII] or dust detection in two Little Red Dots at $z_{\text{spec}} > 7$

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ABSTRACT

Little Red Dots (LRDs) are compact, point-like sources characterized by their red color and broad Balmer lines; it is a matter of debate whether they are dominated by active galactic nuclei (AGNs) or dusty star-forming galaxies (DSFGs). Here we report two LRDs (ID9094 and ID2756) at $z_{\rm spec} > 7$ recently discovered in the JWST FRESCO GOODS-North field. Both satisfy the "v-shaped" color and compactness criteria for LRDs and are identified as Type-I AGN candidates based on their broad H β emission lines (full width at half maximum: $2280 \pm 490 \, \mathrm{km \, s^{-1}}$ for ID9094 and $1070 \pm 240 \, \mathrm{km \, s^{-1}}$ for ID2756) and narrow [O III] lines ($\approx 300 - 400 \, \mathrm{km \, s^{-1}}$). To investigate their nature, we conducted deep NOEMA follow-up observations targeting the [C II] 158 μ m emission line and the 1.3 mm dust continuum. We do not detect [C II] or 1.3 mm continuum emission for either source. If the two LRDs were DSFGs, we would expect significant detections: >16 σ for [C II] and >3 σ for the 1.3 mm continuum of ID9094, and >5 σ for the [C II] of ID2756. Using the 3 σ upper limits of [C II] and 1.3 mm, we performed two analyses: (1) UV-to-far-infrared spectral energy distribution fitting with and without AGN components, and (2) comparison of their properties with the $L_{\rm [C II]}$ -SFR_{tot} empirical relation. Both analyses are consistent with a scenario in which AGN activity contributes to the observed properties, though a dusty star-forming origin cannot be fully ruled out. Our results highlight the importance of far-infrared observations for studying LRDs, a regime that remains largely unexplored.

Key words. galaxies: active - galaxies: high-redshift - galaxies: photometry - galaxies: star formation - submillimeter: galaxies

1. Introduction

Little Red Dots (LRDs; e.g., Fujimoto et al. 2022; Labbe et al. 2025; Matthee et al. 2024; Barro et al. 2024) have emerged as one of the most intriguing populations studied with the JWST (Gardner et al. 2023). The term "LRDs" was first introduced by Matthee et al. (2024) to describe galaxies exhibiting broad $H\alpha$ emission lines with a red point-source morphology. In a broader context, the name LRD has also been applied to a variety of compact, red sources – predominantly at high redshifts ($z \gtrsim 4$) – with "v-shaped" (blue UV and red optical slopes in the rest frame) spectral energy distributions (SEDs), based mainly on photometric measurements only. The samples of these sources overlap, contributing to the complexity of their classification. Most recent studies agree that LRDs involve a combination of stellar and active galactic nucleus (AGN) contributions, although there is ongoing debate over which component dominates (e.g., Wang et al. 2024; Akins et al. 2025a). Some works emphasize a dusty star-forming galaxy-dominated interpretation (e.g., Pérez-González et al. 2024; Labbé et al. 2023; Baggen et al. 2024; Williams et al. 2024), while others find stronger evidence for AGN dominance (e.g., Kocevski et al. 2023, 2025; Greene et al. 2024; Kokorev et al. 2024; Ji et al. 2025; Naidu et al. 2025; de Graaff et al. 2025).

Multiwavelength studies of LRDs have primarily focused on rest-frame UV, optical, near-infrared (NIR), mid-infrared (MIR), X-ray, and radio data, whereas far-infrared (FIR) studies remain relatively scarce in comparison. FIR observations, however, are crucial for disentangling the roles of AGN activity and dusty starburst processes, as they directly probe dust properties, star formation rates (SFRs), and possible AGN contributions. Some studies have started to investigate LRDs in the FIR (e.g., Labbe et al. 2025; Casey et al. 2024, 2025; Akins et al. 2024, 2025b; Williams et al. 2024; Setton et al. 2025), but most rely on photometric redshifts - either through stacked analyses or individual source studies – which might introduce large uncertainties when interpreting their dust properties and star formation activity. Current evidence shows that many LRDs are faint in X-rays (e.g., Ananna et al. 2024; Maiolino et al. 2025; Yue et al. 2024; Sacchi & Bogdan 2025), exhibit weak or no radio emission (e.g., Akins et al. 2024), and lack detectable hot dust in the MIR (e.g., Pérez-González et al. 2024; Williams et al. 2024). These properties could result from heavily obscured AGNs, extreme dusty starbursts, or even hybrid scenarios. Without robust FIR constraints, however, distinguishing between these possibilities remains challenging.

Another potential issue is the partial overlap in photometric selection methods for LRDs (e.g., Labbe et al. 2025; Greene et al. 2024; Kokorev et al. 2024) and optical/NIR-selected dusty star-forming galaxies (DSFGs; optically dark or faint galaxies; e.g., Wang et al. 2019; Franco et al.

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2018; Alcalde Pampliega et al. 2019; Williams et al. 2019, 2024; Gómez-Guijarro et al. 2023, 2022; Xiao et al. 2023; Barrufet et al. 2023; McKinney et al. 2023; Akins et al. 2023; Pérez-González et al. 2023; Barro et al. 2024) due to their similar red colors. This may pose a challenge for DSFG studies, as optical/NIR-selected DSFGs typically have more relaxed selection criteria that rely solely on color and/or magnitude cuts without considering morphology. As a result, LRDs could inadvertently be included in DSFG samples. Additionally, studies also show that some DSFGs have a compact morphology (e.g., Gómez-Guijarro et al. 2023). This overlap introduces ambiguity, as the two populations could appear in the same sample category despite potentially distinct physical origins. For example, DSFGs are typically powered by intense star formation obscured by large amounts of dust, while AGN-driven LRDs may exhibit similar photometric properties due to their reddened SEDs but have fundamentally different drivers of energy output.

In this study we investigated two LRDs, ID9094 and ID2756¹, located in the JWST FRESCO GOODS-North field (Oesch et al. 2023). Due to the overlap in photometric selection criteria between DSFGs and LRDs, these two sources were initially classified as optically faint, dust-obscured galaxy candidates with red colors (F182M – F444W > 1.5 mag) and as the most massive candidates at $z_{\rm spec}$ > 7 (M_{\star} > 10¹⁰ M_{\odot}) based on optical-to-NIR SED fitting (Xiao et al. 2024). However, deep JWST NIRCam grism spectra later revealed the presence of broad H β emission lines (a broad component with a full width at half maximum >1000 km s⁻¹; e.g., Matthee et al. (2024); Sect. 2.4)², along with a red, point-source-like morphology, indicating potential AGN activity.

To further investigate their nature, we present follow-up FIR observations from NOEMA (The Northern Extended Millimetre Array) targeting the [C II] 158 µm emission line and the 1.3 mm dust continuum. This study, together with the recent work by Setton et al. (2025), represents one of the first efforts to investigate LRDs with both dust and [CII] emission. Setton et al. (2025) focus on two of the most luminous LRDs at $z_{\rm spec} \sim 3-4.5$ and report a deficit of hot and cold dust emission. In comparison, our study targets two LRDs at $z_{\rm spec} > 7$, a redshift regime more representative of the typical LRD population identified in JWST surveys. Together, the two studies provide a more complete view of the dust and gas properties of LRDs. Our NOEMA observations provide crucial constraints on the dust temperature, allowing us to assess whether the observed properties of these LRDs are consistent with DSFGs or if AGN-driven processes dominate. By combining FIR data with existing multiwavelength photometry, this work addresses the critical question of whether DSFG scenarios can be ruled out for LRDs in the absence of spectroscopic confirmation.

This paper is organized as follows. In Sect. 2 we describe the data and observational details. Section 3 presents our results, including the non-detection of the dust continuum and [C II] line in the FIR and its implications for the nature of these LRDs. In Sect. 4 we discuss the impact of LRDs on our understanding of DSFGs. We summarize our conclusions in Sect. 5.

Throughout this paper we adopt a Chabrier initial mass function (IMF; Chabrier 2003) to estimate SFRs and stellar masses. We assume a Planck cosmology (Planck Collaboration VI 2020)

with $(\Omega_{\rm m}, \Omega_{\Lambda}, h, \sigma_8) = (0.3, 0.7, 0.7, 0.81)$. When necessary, data from the literature have been converted with a conversion factor of M_{\star} (Salpeter 1955, IMF) = $1.7 \times M_{\star}$ (Chabrier 2003, IMF). All magnitudes are in the AB system (Oke & Gunn 1983), such that $m_{\rm AB} = 23.9 - 2.5 \times \log(S_{\nu} [\mu Jy])$.

2. Data and sample

The two sources ID9094 (α , δ [J2000] = 189.019240, 62.243531) at $z_{\rm spec}$ = 7.0388 ± 0.0001 and ID2756 (α , δ [J2000] = 189.083488, 62.202579) at $z_{\rm spec}$ = 7.1883 ± 0.0001 are located in the GOODS-N field, and were first discovered in the JWST FRESCO survey (Oesch et al. 2023) and reported in Xiao et al. (2024). They have red colors (F182M - F444W > 1.5 mag) and are faint in the F182M band (F182M > 26.4 mag). As such, they were selected as optically faint dusty galaxy candidates, with high inferred stellar masses. In this study, based on their broad H β emission line, "v-shaped" colors, and compact morphology in F444W, we now identify both sources as LRDs, following commonly used criteria (e.g., Matthee et al. 2024; Labbe et al. 2025; Greene et al. 2024; Kokorev et al. 2024; Kocevski et al. 2025; see details in Sect. 2.4).

2.1. JWST observations

The two sources were initially observed by the JWST FRESCO survey (Oesch et al. 2023), which provides NIRCam/grism spectroscopy in the F444W filter and direct imaging in F182M, F210M, and F444W over ~62 arcmin² in each GOODS field, North and South. The grism spectra cover a wavelength range of 3.8–5.0 μm at a resolution of R \sim 1600, reaching a 5σ line sensitivity of $\sim 1.3 \times 10^{-18}$ erg s $^{-1}$ cm $^{-2}$. The images reach typical 5σ depths of 28.3, 28.1, and 28.2 mag, respectively, in a 0.16 circular aperture. Data reduction was performed using the grizli pipeline³, including continuum subtraction and optimal 1D spectral extraction.

Recently, we obtained deeper NIRCam/grism spectroscopy from a Cycle 3 JWST program (GO-4762; PIs: S. Fujimoto & G. Brammer), using the F410M filter instead of F444W. This choice improves spectral sensitivity by blocking the high thermal background beyond 4.4 μm , at the cost of reduced wavelength coverage. Despite a slightly shorter exposure time (by a factor of 1.3) than FRESCO, this dataset thus achieves a 5σ line sensitivity of $\sim\!7.8\times10^{-19}\,\mathrm{erg\,s^{-1}\,cm^{-2}}$. The continuum subtraction was optimized to explicitly mask emission lines and a wide kernel with a large central hole (151 and 31 pixels, respectively; for the running median) to prevent the over-subtraction of broad wings (see Matthee et al. 2024).

The NIRCam images and grism spectra of ID9094 and ID2756 are shown in Fig. 1, updated with the latest data.

2.2. NOEMA observations

The NOEMA observations of ID9094 and ID2756 were carried out in the summer of 2023 using the PolyFiX correlator (Project ID: S23CY, PIs: D. Elbaz & M. Xiao). Conducted with array configuration D in band 3, the observations targeted the [C II] 158 μm emission line, and the 1.3 mm dust continuum, with representative frequencies of 232.052 GHz for ID9094 and 236.576 GHz for ID2756. The total on-source time was 6 hours and 4.5 hours for ID2756 and ID9094, respectively.

 $^{^1\,}$ ID9094 and ID2756 correspond to gn17537 and gn28924, respectively, in the FRESCO [O III] catalog (Meyer et al. 2024).

² The broad lines were not clearly identified in the initial FRESCO spectra or at the time of the NOEMA proposal, due to the limited sensitivity of the FRESCO data.

³ https://github.com/gbrammer/grizli

Table 1. Flux or flux densities of two LRDs from multiwavelength observations.

Instrument	$\lambda_{ m ref}$	Units	ID2756	ID9094	
Chandra ^a	0.5-7 keV	${\rm erg~s^{-1}~cm^{-2}}$	$< 1.1 \times 10^{-16}$	$<1.1 \times 10^{-16}$	
	2-7 keV	${\rm erg~s^{-1}~cm^{-2}}$	$< 1.8 \times 10^{-16}$	$<1.8 \times 10^{-16}$	
	$0.5-2 \mathrm{keV}$	${\rm erg~s^{-1}~cm^{-2}}$	$< 3.6 \times 10^{-17}$	$< 3.6 \times 10^{-17}$	
HST/ACS F435W	$0.43 \mu m$	nJy	<19.4	<17.7	
HST/ACS F606W	0.59 µm	nJy	<12.3	< 9.4	
HST/ACS F775W	0.77 μm	nJy	<13.6	<17.7	
HST/ACS F814W	0.80 µm	nJy	<17.1	<13.5	
HST/ACS F850LP	0.90 µm	nJy	<19.5	<24.4	
HST/WFC3 F105W	1.06 µm	nJy 33.0 ± 8.4			
HST/WFC3 F125W	1.25 μm	nJy	87.5 ± 10.8	75.7 ± 9.8	
HST/WFC3 F140W	$1.40\mu m$	nJy	84.7 ± 6.6		
HST/WFC3 F160W	1.54 µm	nJy	91.3 ± 7.8	109.7 ± 10.4	
JWST/NIRCam F090W	$0.90\mu m$	nJy	<14.3	<21.0	
JWST/NIRCam F115W	1.15 µm	nJy	79.2 ± 4.4	64.8 ± 6.2	
JWST/NIRCam F150W	1.50 µm	nJy		93.7 ± 12.9	
JWST/NIRCam F182M	1.85 µm	nJy	80.0 ± 4.0	99.3 ± 6.7	
JWST/NIRCam F200W	1.99 µm	nJy		109.6 ± 11.5	
JWST/NIRCam F210M	2.10 µm	nJy	117.8 ± 5.9	130.2 ± 8.8	
JWST/NIRCam F356W	$3.57 \mu m$	nJy	225.8 ± 11.3	443.8 ± 22.2	
JWST/NIRCam F444W	$4.40\mu m$	nJy	372.4 ± 18.6	887.7 ± 44.4	
Spitzer/IRAC	$3.6\mu m$	nJy	191.0 ± 15.4	400.4 ± 13.0	
	$4.5 \mu m$	nJy	340.2 ± 17.6	592.7 ± 16.5	
	5.8 µm	nJy	933.7 ± 243.5	<676.5	
	$8.0\mu m$	nJy	351.2 ± 284.8	2683.0 ± 289	
Spitzer/MIPS ^b	$24 \mu m$	μ Jy	<21	<21	
	70 μm	mJy	<2.4	<2.4	
Herschel/PACS ^b	$100 \mu m$	mJy	<1.1	<1.1	
	160 μm	mJy	<2.7	<2.7	
$Herschel/SPIRE^b$	$250\mu m$	mJy	< 5.7	< 5.7	
	$350\mu m$	mJy	<7.2	<7.2	
	$500 \mu m$	mJy	<9	<9	
SCUBA-2 ^c	$850\mu m$	mJy	<1.2	<1.2	
NOEMA	1.3 mm	μ Jy	<151	<91	
	$[C{\rm II}]_{158\mu m}$	$\mathrm{Jy}\mathrm{km}\mathrm{s}^{-1}$	< 0.310	< 0.157	
$NIKA-2^d$	1.2 mm	μ Jy	< 510	< 510	
	2.0 mm	μ Jy	<144	<144	
VLA 1.5 GHz ^e	20 cm	uJy	<6.6	< 6.6	

Notes. We use 1σ for uncertainties and 3σ upper limits for non-detections. (a) We calculate 3σ upper limits based on the mean sensitivity limits achieved in the central ~1 arcmin² area at the average aim point from 2 Ms *Chandra* Deep Field-North survey (Xue et al. 2016). (b) Sources are not detected, we adopt 3σ upper limits from the GOODS-Herschel survey (Elbaz et al. 2011). (c) We adopt the confusion limit for a 3σ detection in the central region of GOODS-N from Cowie et al. (2017). (d) We adopt 3σ upper limits according to the N2CLS depth from Bing et al. (2023). (e) Calculated based on the rms noise for a 1"6 resolution in the field center from Owen (2018).

All the calibrations and the creation of *uv* table were performed using CLIC package from the IRAM GILDAS software⁴ with the support of IRAM astronomers. We then performed further analysis with version 6.5.5 of the Common Astronomy Software Application (CASA) package (McMullin et al. 2007).

For the [C II] 158 μ m emission line, imaging was carried out using the tclean task with 0.4" pixels and a channel width of 50 km s⁻¹ with a natural weighting. The natural weighting provides the best point-source sensitivity, which is optimal for source detection. The resulting data cube has a synthesized beam size of full width at half maximum (FWHM) $\simeq 1".838 \times 1".472$ (~ 9.8 kpc $\times 7.8$ kpc in physical scale) and $2".859 \times 1".854$ (~ 15.0 kpc $\times 9.7$ kpc) with a root mean square (rms) sensitivity of ~ 0.45 and 0.84 mJy beam⁻¹ per channel at the phase center, for ID9094 and ID2756, respectively.

We also created the observed 1.3 mm continuum maps using natural weighting and the same method as for the [C II] 158 μ m emission line. For ID9094, the rms level is ~30.4 μ Jy beam⁻¹ in the map of 1".725 × 1".383 (~9.2 kpc × 7.4 kpc) angular resolution. For ID2756, the rms level is ~50.3 μ Jy beam⁻¹ in the map of 2".712 × 1".743 (~14.3 kpc × 9.2 kpc) angular resolution.

2.3. Multiwavelength observations

Besides the JWST FRESCO and NOEMA observations, the two sources are also covered by a wide array of multiwavelength observations. In Table 1, we summarize the multiwavelength dataset for our two sources from X-ray to radio: (i) X-ray: Chandra 2 Ms (0.5-7.0 keV, 0.5-2.0 keV, and 2-7 keV bands) images in the Chandra Deep Field-North field (CDF-N; Xue et al. 2016); (ii) Hubble Space Telescope (HST) ACS (F435W, F606W, F775W, F814W, F850LP) and HST/WFC3 (F105W, F125W, F140W, F160W) images from the *Hubble* Legacy Fields (HLF)⁵ program (Whitaker et al. 2019); (*iii*) JWST/NIRCam (F090W, F115W, F150W, F182M, F200W, F210M, F356W, and F444W) from FRESCO (Oesch et al. 2023), CONGRESS (Egami et al. 2023), and JADES (Eisenstein et al. 2023; D'Eugenio et al. 2025). The two sources are not yet covered by any publicly available JWST/MIRI observations; (iv) Spitzer/IRAC from the GREATS program (3.6 µm, 4.5 µm, 5.8 µm, 8 µm; Stefanon et al. 2021, new catalogs based on FRESCO sources are taken from Stefanon et al. in prep.; priv. comm.); (v) Spitzer/MIPS, Herschel/PACS, and SPIRE catalog is taken from the GOODS-Herschel survey (Elbaz et al. 2011); (vi) JCMT/SCUBA-2 at 850 μm (Cowie et al. 2017); (vii) NOEMA data from this study (see Sect. 2.2); (viii) NIKA2 Cosmological Legacy Survey at 1.2 mm and 2.0 mm (Bing et al. 2023); and (ix) VLA 1.5 GHz (Owen 2018).

We note that the photometric measurements from HST and JWST images are derived following the same procedure as outlined in Weibel et al. (2024). Briefly, we used SExtractor (Bertin & Arnouts 1996) in dual image mode with an inverse-variance weighted stack of the F210M + F444W with FRESCO-only imaging, as the detection image. In this study, fluxes were measured in 0".16 radius circular apertures in images that are point spread function (PSF)-matched to the F444W band. Total fluxes were derived from the Kron AUTO aperture provided by SExtractor in the F444W band, in addition to a correction based on the encircled energy of the Kron aperture on the F444W PSF. Detailed descriptions of data reduction and photometric measurements are provided in Weibel et al. (2024) and Xiao et al. (2024).

2.4. Our sample: Two "Little Red Dots" at $z_{\rm spec} > 7$

Although there is no uniform definition of LRDs, we identified our two sources as LRDs because they satisfy both of two commonly used selection criteria: (1) an initial criterion – sources with broad Balmer emission lines ($\nu_{\rm FWHM,broad} > 1000 \, \rm km \, s^{-1}$; Matthee et al. 2024); and (2) a widely used criterion based only on photometric measurements – a combination of "v-shaped" SED and compactness selection (e.g., Labbe et al. 2025; Kokorev et al. 2024; Greene et al. 2024; Kocevski et al. 2025).

Figure 1 presents the JWST/NIRCam images and grism spectra of ID9094 and ID2756. The red-green-blue (RGB) images show that both sources appear red with point-like mor-

⁴ https://www.iram.fr/IRAMFR/GILDAS/

⁵ https://archive.stsci.edu/prepds/hlf/

Table 2. Best-fit emission line properties from JWST/NIRCam F410M grism spectra, and best-fit UV and optical slopes from photometric bands.

ID	$z_{ m spec}$	Flux _{Hβ} ,broad 10^{-18} erg s ⁻¹ cm ⁻²	Flux _{Hβ,narrow} 10^{-18} erg s ⁻¹ cm ⁻²	Flux _[O III] , 15008 10 ⁻¹⁸ erg s ⁻¹ cm ⁻²	FWHM _{Hβ,broad} km s ⁻¹	FWHM _{narrow} km s ⁻¹	$eta_{ m UV}$	$eta_{ m opt}$
, , , .	=	6.95 ± 1.57 2.65 ± 0.65	2.01 ± 0.63 0.52 ± 0.40	24.47 ± 0.49 6.69 ± 0.29	2280 ± 490 1070 ± 240	425 ± 10 295 ± 16	-1.05 ± 0.24 -1.48 ± 0.12	

Notes. β_{UV} and β_{opt} are continuum slopes derived from multi-band JWST photometry blueward and redward of the Balmer break, respectively, and trace the two sides of the "v-shaped" SEDs (see Sect. 2.4).

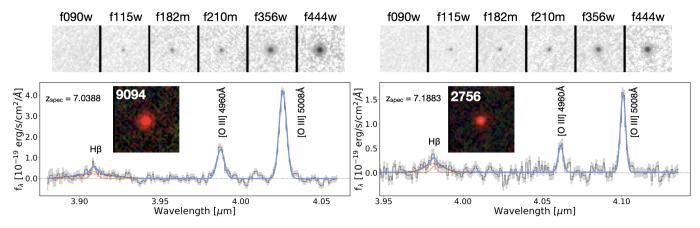


Fig. 1. JWST images and spectra of the two sources. Top: $2'' \times 2''$ stamps obtained in JWST/NIRCam filters (0.90 μm, 1.15 μm, 1.82 μm, 2.10 μm, 3.56 μm, and 4.44 μm). Bottom: 1D spectra (covering H β , [O III] $\lambda\lambda$ 4960, 5008 Å emission lines) obtained from NIRCam/grism observations with the F410M filter, with RGB images embedded (F182M in blue, F210M in green, and F444W in red). The gray-shaded areas show the associated 1σ uncertainty. The best-fit Gaussian line model is shown in blue. Both sources have a broad H β line; the solid orange line shows the narrow component, and the dashed red line shows the broad component.

phology. The grism spectra show H β and [O III] $\lambda\lambda$ 4960, 5008 emission lines. Both sources exhibit a broad H β line. We fit the H β line with both narrow and broad Gaussian components, assuming the narrow component has the same width as the [O III] $\lambda\lambda$ 4960, 5008 lines. The best-fit line properties are shown in Table 2. The broad component of H β has a FWHM of $\nu_{\text{FWHM,H}\beta,\text{broad}} = 2280 \pm 490 \,\text{km s}^{-1}$ for ID9094 and $\nu_{\text{FWHM,H}\beta,\text{broad}} = 1070 \pm 240 \,\text{km s}^{-1}$ for ID2756, suggesting a possible Type I (broad emission line) AGN. Notably, ID2756 has also been identified as a strong Lyman- α emitter, with a considerable equivalent width of 221 Å (Leonova et al. in prep.).

We further compared our two sources to an alternative color and compactness selection criterion for LRDs, designed to identify compact, red objects with "v-shaped" SEDs (e.g., Kokorev et al. 2024; Greene et al. 2024):

$$\begin{split} & \text{S/N(F444W)} > 14 \quad \& \quad \text{m}_{\text{F444W}} < 27.7 \, \text{mag}, \\ & -0.5 < \text{F115W} - \text{F200W} < 1.0, \\ & \text{F277W} - \text{F444W} > 1.0, \\ & \text{F277W} - \text{F356W} > 0.7, \\ & \text{compact} = f_{\text{F444W}}(0\rlap.''4)/f_{\text{F444W}}(0\rlap.''2) < 1.7. \end{split}$$

Both sources are bright in F444W, with magnitudes of 24.0 ± 0.1 mag (ID9094) and 25.0 ± 0.1 mag (ID2756), consistent with the LRD selection. Since ID2756 lacks observations in F277W and F200W, while ID9094 is missing F277W, the corresponding fluxes used here are derived from the best-fit SED and the 16th–84th percentile range (as shown later in Fig. 4). The detailed comparison is shown in Fig. 2. ID9094 lies within the red region, perfectly matching the LRD selection. ID2756 exhibits a "v-shaped" color and compact morphology, aligning with the LRD classification. Its F277W – F356W color is slightly bluer than the

selection threshold but still within the 1σ uncertainty. We therefore classify both sources as LRDs. Compared to the majority of sources from various JWST blank legacy fields (blue-shaded region in Fig. 2), obtained from the DAWN JWST Archive (DJA 6 ; Valentino et al. 2023), these two LRDs stand out as particularly exceptional, exhibiting redder colors and more compact morphologies.

To further quantify the "v-shaped" SEDs, we also measured the UV and optical continuum slopes ($\beta_{\rm UV}$ and $\beta_{\rm opt}$) using multi-band photometry, following the method of Kocevski et al. (2025). Both sources satisfy their LRD selection criteria (-2.8 < $\beta_{\rm UV}$ < -0.37 and $\beta_{\rm opt}$ > 0). The measured slope values are listed in Table 2. In addition, both sources are unresolved in F444W, further supporting that they are typical LRDs.

Altogether, we find that both sources are consistent with standard criteria for LRDs based on their morphology and "v-shaped" SED. In addition, both sources show the broad H β line, making them strong Type I AGN candidates.

3. Results

3.1. Non-detections in both [C II] 158 μm and 1.3 mm of two LRDs

We analyzed the NOEMA data cubes and extracted spectra at the positions of our sources using an aperture consistent with the beam size. The resulting spectra, shown in Fig. 3, reveal no significant signal at the expected [C II] frequency ($\ll 3\sigma$). Assuming a [C II] 158 µm line width of 250 km s⁻¹ – a typical value for high-redshift galaxies (Béthermin et al. 2020; Endsley et al. 2022) – we constructed the velocity-integrated intensity map

⁶ https://dawn-cph.github.io/dja/imaging/v7/

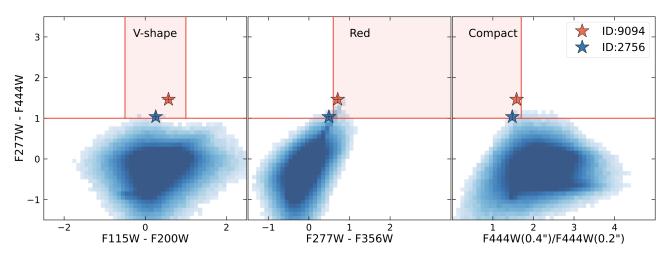


Fig. 2. Locations of the two sources (red and blue stars) relative to the typical color and compactness selection criteria for LRDs. The red-shaded regions highlight LRDs that have a "v-shaped" SED (*left*), red colors (*middle*), and a compact morphology (*right*), defined by the criteria: -0.5 < F115W - F200W < 1.0, F277W - F444W > 1.0, F277W - F356W > 0.6, and $f_{F444W}(0.4)/f_{F444W}(0.4)/f_{F444W}(0.4)/f_{F444W}(0.4)/f_{F444W}(0.4)$. Error bars indicate 1σ uncertainties. The blue-shaded region represents the distribution of sources observed in these filters from the prime JWST blank legacy fields in the DIA.

(moment-0). Figure 3 overlays the contours from the [C II] moment-0 and 1.3 mm maps on the JWST RGB image, showing no detections ($\ll 3\sigma$) at the source positions and in the surrounding regions. We therefore conclude that neither the [C II] 158 μ m emission line nor the 1.3 mm dust continuum emission are detected for the two LRDs in the current NOEMA data.

We then determined 3σ upper limits for both [C II] and dust continuum. Using the rms of the 1.3 mm map (see Sect. 2.2) and assuming our sources are not resolved, we derive 3σ upper limits of 91 μ Jy and 151 μ Jy for ID9094 and ID2756, respectively. In addition, the 3σ upper limits for [C II] are calculated as three times the pixel-by-pixel rms of the moment-0 map, which gives upper limits of 0.157 Jy km s⁻¹ and 0.310 Jy km s⁻¹ for ID9094 and ID2756, respectively. The corresponding 3σ upper limits on the [C II] emission-line luminosity are $\log(L_{\rm [C\,II]}/L_{\odot}) < 8.23$ for ID9094 and $\log(L_{\rm [C\,II]}/L_{\odot}) < 8.58$ for ID2756. Using $L_{\rm [C\,II]}$, we further calculate the 3σ upper limits on the molecular gas mass (Zanella et al. 2018) to be $\log(M_{\rm mol}/M_{\odot}) < 9.7$ for ID9094 and $\log(M_{\rm mol}/M_{\odot}) < 10.1$ for ID2756.

Intriguingly, when requesting the NOEMA observations, we assumed these two galaxies to be DSFGs. Given their stellar masses from Xiao et al. (2024), we estimated their total SFRs (SFR_{tot}) assuming that they lie on the star-forming main sequence (Schreiber et al. 2015) and predicted their expected [C II] fluxes using the $L_{\text{[C II]}}$ -SFR_{tot} relation for SFGs (Schaerer et al. 2020). Similarly, we estimated their 1.3 mm continuum fluxes by assuming an infrared SED shape consistent with typical dust-obscured galaxies (optically dark or faint galaxies; Xiao et al. 2023). Given the depth of our NOEMA observations, we expected ID9094 to be detected at >16 σ in [C II] and $>3\sigma$ in the 1.3 mm continuum, while ID2756 should be detected in [C II] at $>5\sigma$. However, the complete absence of [C II] and 1.3 mm emission suggests that these two LRDs are not the typical DSFGs initially assumed. This result aligns with our findings in Sect. 2.4: both sources exhibit a broad H β emission line, being Type-I AGN candidates.

3.2. The hints of AGNs in two LRDs from FIR non-detections

In this section we utilize the 3σ upper limits on the [C II] 158 μm emission line and the 1.3 mm continuum to investigate the pos-

sible presence of AGN activity in the LRDs. While both sources have already been identified as strong Type-I AGN candidates based on their broad H β emission lines, our primary goal here is to investigate whether the FIR non-detections can effectively rule out the DSFG scenario for LRDs, independently of spectral information. Here, we focus primarily on ID9094, which benefits from deeper NOEMA observations compared to ID2756, allowing for more robust constraints. We assessed the AGN contribution through two independent approaches: (1) UV-to-FIR SED fitting analysis and (2) the $L_{\rm [CII]}$ -SFR_{tot} empirical relation.

3.2.1. SED fitting analysis

We performed a UV-to-FIR SED fitting analysis using CIGALE (Boquien et al. 2019), which is based on an energy balance principle, fixing the redshift to $z_{\rm spec}$. We assumed a constant star formation history and adopted Bruzual & Charlot (2003) stellar population models with the Calzetti et al. (2000) dust attenuation law. We adopted a broad metallicity range of 0.004-0.02, dust attenuation to the rest-frame V band (A_V) values from 0to 6 magnitudes, and ionization parameters $\log U$ of -2.0. The AGN component was modeled using the Fritz et al. (2006) template, with the AGN fraction (f_{AGN}) , i.e., the contribution of the AGN to the total IR luminosity, ranging from 0.01 to 0.8. The angle between the equatorial axis and the viewing angle at the line of sight (Ψ) ranges from 40 to 90. With a fixed open angle of the AGN torus at 100 degrees, this range of Ψ refers to Type I AGNs. We note that recent MIR observations (Wang et al. 2025; Williams et al. 2024; Barro et al. 2024) suggest a lack of hot torus dust emission in LRDs, which may deviate from the Fritz et al. (2006) template adopted. However, with the absence of deep MIR observations, we anticipate little change in the quality of SED fitting on our sample by only changing the AGN torus emission in the templates. On the contrary, the inclusion of continuum from accretion disk by limiting Ψ to Type I AGNs is critical, as we robustly detected broad lines in H β .

Given that we only have upper limits in the FIR, we adopted a simple FIR dust model from Casey (2012), using the default dust emissivity index ($\beta = 2.0$) and MIR power-law slope ($\alpha = 2.0$). Compared to $\beta = 1.5$, which is adopted in some studies of DSFGs (e.g., Hildebrand 1983; Kovács et al. 2006; Gordon et al.

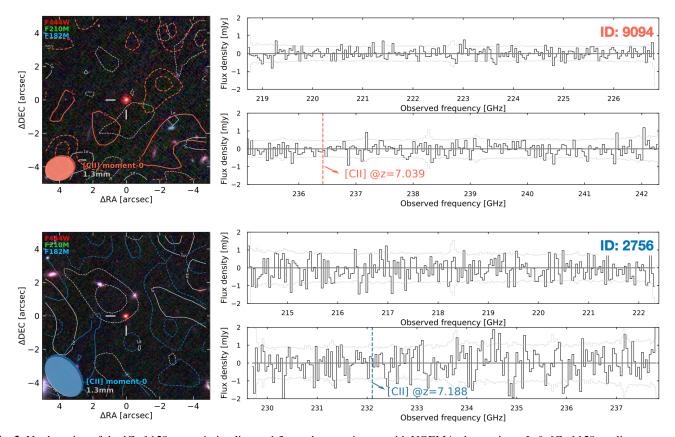


Fig. 3. No detection of the [C II] 158 μ m emission line or 1.3 mm dust continuum with NOEMA observations. *Left*: [C II] 158 μ m line moment-0 map (assuming a line width of 250 km s⁻¹) and 1.3 mm dust contours overlaid on the JWST RGB image (10" × 10"). The contour levels start at 1σ and increase in steps of $\pm 1\sigma$, where positive and negative contours are solid and dashed, respectively. The beam sizes are displayed in the lower-left corner. *Right*: NOEMA 1.3 mm spectra with 50 km s⁻¹ binning. The gray lines show the associated 1σ uncertainty. The vertical dashed lines highlight the locations of [C II] 158 μ m. The spectra are taken at the positions of the sources and match the apertures of the beam sizes, assuming the emission line is unresolved.

2010; Gómez-Guijarro et al. 2022; Xiao et al. 2023), our choice of $\beta=2.0$ is more conservative and leads to slightly lower inferred dust temperatures for a given FIR slope. During the fitting process, we find that when the dust temperature $(T_{\rm dust})$ is allowed to vary across a defined range, the best-fit $T_{\rm dust}$ is strongly influenced by the upper boundary of the input range. This effect arises due to the absence of deep MIR data, which prevents us from constraining the maximum $T_{\rm dust}$. Consequently, we cannot determine a precise best-fit value for $T_{\rm dust}$, but the deep NOEMA data allow us to place a lower limit.

To refine our constraints, we adopted an optimized fitting approach that allowed us to constrain the minimum allowed T_{dust} using the NOEMA 1.3 mm 3σ upper limits. In this approach, we fixed T_{dust} during the fitting and performed multiple fits with different $T_{\rm dust}$ values, ranging from 30 K to 150 K. In addition, we also corrected the effect of the cosmic microwave background (CMB) as an observing background on photometry, following da Cunha et al. (2013). The CMB has the effect of raising the apparent dust temperature with increasing redshift, as the CMB temperature scales with redshift and provides an additional background source of radiation for dust heating. We assumed $\beta = 2.0$, consistent with previous studies for $z \sim 7$ galaxies (Sommovigo et al. 2022). In this way, we derived the intrinsic FIR photometry starting from MIPS 70 µm, which was used for SED fitting. The best-fit SEDs, with and without AGN components considered, are presented in Fig. 4.

For ID9094, without including an AGN component (Fig. 4-top), we present three representative examples of best-fit SEDs

with (a) $T_{\text{dust}} = 40 \text{ K}$, (b) $T_{\text{dust}} = 110 \text{ K}$, and (c) $T_{\text{dust}} =$ 150 K. These correspond to results that the predicted flux density at 1.3 mm exceeds, matches, and falls below the observed 3σ upper limits, respectively, under the energy balance assumption of CIGALE. At $T_{\text{dust}} = 40 \,\text{K} - \text{a typical value for star-}$ forming UV-bright galaxies in the REBELS survey at $z \sim 7$ (Sommovigo et al. 2022) – the predicted 1.3 mm flux exceeds the observed 3σ upper limit by nearly a factor of 3. This discrepancy suggests that, with our current NOEMA observations, we should have detected the source at a significance level of $\sim 8\sigma$, which is clearly not the case. As $T_{\rm dust}$ increases, the peak of the FIR SED moves toward the shorter wavelength and the fit improves. At $T_{\text{dust}} = 110 \,\text{K}$ (Fig. 4-b), the predicted 1.3 mm flux aligns with the observed 3σ upper limit. However, given that the actual signal-to-noise ratio at 1.3 mm is well below 3σ , this implies that the real $T_{\rm dust}$ would be greater than 110 K.

Therefore, without including an AGN component, we obtain a lower limit of $T_{\rm dust} \gtrsim 110\,\rm K$. At fixed $T_{\rm dust} = 110\,\rm K$, as in Table 3, the best-fit IR luminosity is $L_{\rm IR} = 13.8^{+1.8}_{-2.6} \times 10^{11}\,L_{\odot}$, SFR = $125.2^{+8.7}_{-8.6}\,\rm M_{\odot}\,\rm yr^{-1}$ (we use this for the analysis in Sect. 3.2.2), and stellar mass of $M_{\star} = 2.3^{+0.2}_{-0.2} \times 10^{10}\,M_{\odot}$. Remarkably, even this lower limit is significantly higher than the typical $T_{\rm dust} \sim 40\,\rm K$ observed in REBELS galaxies at similar redshifts (Sommovigo et al. 2022). Similar findings have been reported in studies of 675 LRDs at $z \gtrsim 4$, regardless of heating mechanism, whether AGNs or star formation, where an average luminosity-weighted dust temperature of $T_{\rm dust} = 110\,\rm K$ is reported (Casey et al. 2024). Such an extreme dust temperature

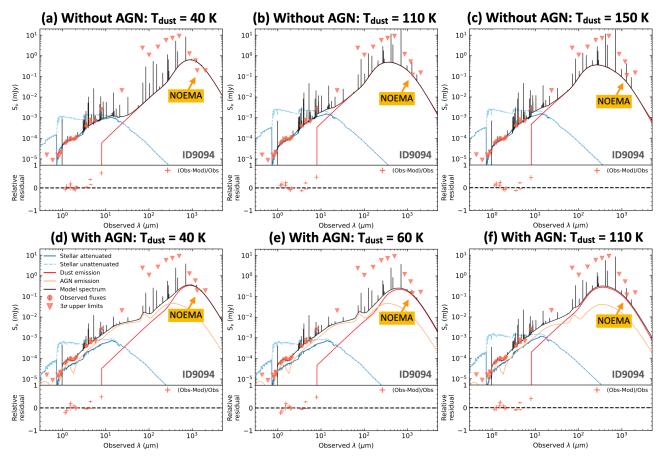


Fig. 4. Best-fit SEDs of ID9094 with fixed dust temperatures, providing the first hint of an AGN scenario in this work. *Top panels*: Best-fit SEDs without an AGN component. *Bottom panels*: Best-fit SEDs with an AGN component. The 3σ upper limit from NOEMA 1.3 mm, the deepest FIR photometric constraint, is highlighted; it plays a key role in shaping the FIR SED and constraining the lower limit of T_{dust} . The FIR photometry has been corrected for CMB effects. We show three representative T_{dust} during the fitting with and without an AGN component, which results in the predicted 1.3 mm flux exceeding (a and d), equal to (b and e), and below (c and f) the observed 3σ upper limit. Panel d shows the fit with an AGN and $T_{\text{dust}} = 40 \text{ K}$, which slightly exceeds the 3σ limit but is not ruled out given the model degeneracy and observational uncertainties (see Sect. 3.2.1). Without an AGN, the SED fitting requires $T_{\text{dust}} \gtrsim 110 \text{ K}$. This value significantly exceeds the typical $T_{\text{dust}} \sim 40 \text{ K}$ observed in REBELS galaxies at similar redshifts (Sommovigo et al. 2022), and may indicate additional dust heating mechanisms such as AGN activity. In contrast, including an AGN component (bottom panels) allows acceptable fits with T_{dust} as low as \sim 40–60 K, similar to AGNs and/or quasar host galaxies at z > 4 (\sim 40–100 K; e.g., Walter et al. 2022; Tsukui et al. 2023; Tripodi et al. 2023; Decarli et al. 2023), supporting the presence of an AGN in ID9094.

Table 3. CIGALE best-fit results for ID9094, with and without an AGN component.

ID9094	fixed T_{dust} [K]	$\log M_{\star} [M_{\odot}]$	SFR [M_{\odot} yr ⁻¹]	$L_{ m IR}\left[L_{\odot} ight]$
Without AGN	110	$2.3^{+0.2}_{-0.2} \times 10^{10}$	$125.2^{+8.7}_{-8.6}$	$13.8^{+1.8}_{-2.6} \times 10^{11}$
With AGN	60	$9.7^{+3.0}_{-8.6} \times 10^9$	$33.1^{+11.4}_{-12.9}$	$6.2^{+2.2}_{-4.3} \times 10^{11}$
With AGN	40	$4.1^{+5.8}_{-2.9} \times 10^9$	$22.4_{-3.9}^{+10.8}$	$5.0^{+3.2}_{-1.6} \times 10^{11}$

is challenging to explain solely through heating by star formation. In typical star-forming galaxies, the dust temperature is primarily driven by the ultraviolet radiation field of young stars, which is usually insufficient to reach the high $T_{\rm dust}$ values observed here (see, e.g., Harshan et al. 2024; Laporte et al. 2017).

Based on the best-fit $L_{\rm IR}$ without AGNs at fixed $T_{\rm dust} = 110\,\rm K$, we derived the dust emitting region size $R_{\rm d} \sim 75\,\rm pc$, based on a blackbody assumption⁷. A crucial parameter to consider is the stellar surface density, given by $\Sigma_{\star} = M_{\star}/2\pi R_{\rm d}^2 \sim 1000$

 $10^6\,M_\odot\,\mathrm{pc}^{-2}$ (similar findings are reported in, e.g., Baggen et al. 2024). This is about one order of magnitude higher than the densest systems in the local universe (Hopkins et al. 2010) and star-forming galaxies at $z\sim5$ –14 observed by JWST (Morishita et al. 2024; Schaerer et al. 2024), highlighting a general issue for LRD models that rely solely on stellar components. Expanding the star formation region would reduce Σ_\star , but it would also decrease $T_{\rm dust}$ below the observed values (e.g., De Rossi et al. 2018), leading to an inconsistency with our findings. These values are difficult to explain through star formation alone under standard assumptions and may suggest the contribution of additional heating sources, such as AGN activity.

⁷ $L_{\rm IR} \sim 4\pi R_{\rm d}^2 \sigma T_{\rm dust}^4$ and σ is the Stefan–Boltzmann constant.

However, we emphasize that this interpretation remains model-dependent, and we further discuss such caveats in Sect. 3.2.3.

When the AGN component is included in the SED fitting (Fig. 4-bottom), the results become more physically plausible. Similar to Fig. 4-top, we also show three representative examples of best-fit SEDs with (d) $T_{\text{dust}} = 40 \text{ K}$, (e) $T_{\text{dust}} = 60 \text{ K}$, and (f) $T_{\text{dust}} = 110 \,\text{K}$. However, due to the degeneracy between the dust temperature and the AGN fraction, the current upper limit at 1.3 mm does not allow us to place a meaningful lower limit on $T_{\rm dust}$. The observed photometry is consistent with SED fits with $T_{\rm dust} = 60 \,\rm K$ and even marginally consistent with $T_{\rm dust}$ down to ~40 K, a value comparable to those observed in AGNs and/or quasar host galaxies at z > 4 ($\sim 40-100$ K; e.g., Walter et al. 2022; Tsukui et al. 2023; Tripodi et al. 2023; Decarli et al. 2023; Meyer et al. 2025). At fixed $T_{\text{dust}} = 60 (40) \text{ K}$, the best-fit AGN fraction is $f_{AGN} = 0.44 \pm 0.12 (0.58 \pm 0.12)$, indicating a scenario where both AGNs and star formation contribute significantly to the FIR emission. The corresponding best-fit $L_{\rm IR}$, SFR, and M_{\star} are listed in Table 3. We also note that both sources are covered by the MEOW survey (JWST Cycle 3 GO-5407; PI: G. Leung; priv. comm.), which detects both in F1000W and ID9094 as well as in F2100W. Although the fluxes are not yet publicly released, preliminary analysis suggests that ID9094 exhibits significant hot dust emission in the MIR (G. Leung, priv. comm.).

For ID2756, we performed SED fitting using the same strategy as for ID9094 (see Appendix A). However, the current 1.3 mm observations for ID2756 (rms level $\sim\!50.3\,\mu\mathrm{Jy\,beam^{-1}})$ are not sufficiently deep to constrain the SED effectively. The lack of sensitivity prevents us from distinguishing between the DSFG and AGN scenarios based solely on the SED fitting results. To achieve meaningful constraints, deeper FIR/submillimeter observations are required. Specifically, observations at least as deep as for ID9094 (rms level $\sim\!30.4\,\mu\mathrm{Jy\,beam^{-1}})$ are necessary to provide tighter limits on the 1.3 mm flux density and help refine the SED fitting.

3.2.2. L_{IC} -SFR_{tot} relation

In this section we compare the locations of our LRDs with the well-established $L_{\rm [CII]}$ –SFR $_{\rm tot}$ scaling relation to investigate their alignment with expectations for star-forming galaxies or the potential influence of AGN activity. The $L_{\rm [CII]}$ –SFR $_{\rm tot}$ relation has been shown to hold across a wide range of redshifts, with no significant evolution observed up to $z\sim 9$ (e.g., Lagache et al. 2018; Schaerer et al. 2020; Carniani et al. 2020; Pallottini et al. 2022). Additionally, studies indicate that the relation is consistent for both AGN-hosting galaxies and typical star-forming galaxies, suggesting that the presence of AGNs does not significantly alter the global $L_{\rm [CII]}$ –SFR $_{\rm tot}$ relation (e.g., Herrera-Camus et al. 2018).

For ID9094, when no AGN component is included in the SED fitting, the SFR derived corresponds to $T_{\rm dust}=110~\rm K$. Combining this SFR with the 3σ upper limits of the $L_{\rm [C\,II]}$ ($\log(L_{\rm [C\,II]}/L_{\odot})<8.23$; Sect. 3.1) places ID9094 significantly below the $L_{\rm [C\,II]}$ –SFR tot relation, with a deviation of \sim 1 dex (see Fig. 5). This large offset suggests that the observed properties of ID9094 may not be fully explained by typical star formation processes in DSFGs. In contrast, when including an AGN component in the SED fitting, the resulting SFR with the $T_{\rm dust}=60~\rm K$ template is about four times lower than that returned by SED fitting without an AGN component. In this case, ID9094 lies within the scatter of the $L_{\rm [C\,II]}$ –SFR tot relation, indicating consistency with the scaling relation. We note that including an AGN compo-

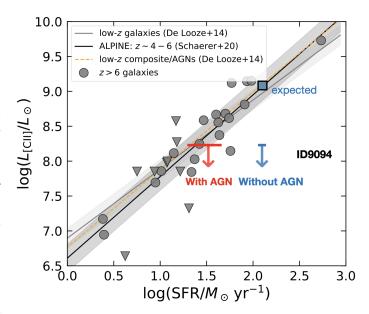


Fig. 5. [CII] luminosity as a function of SFR, providing the second hint of an AGN scenario in this work. ID9094 is shown with 3σ upper limits on $L_{\rm [CII]}$ ($L_{\rm [CII]}$ < 1.7 × 10⁸ L_{\odot}) in red and blue arrows, representing cases with (fixed $T_{\rm dust} = 60 \, \rm K$) and without (fixed $T_{\rm dust} = 110 \, \rm K$) an AGN component, respectively. We note that results for $T_{\rm dust} = 40 \, \rm K$ with an AGN component are not shown here, as they yield similar or even better agreement with the relation (see Sect. 3.2.2). The blue square indicates the expected $L_{\rm [C\,II]}$ location of ID9094 if it is a starforming galaxy without an AGN. Solid gray and black lines, along with their shaded regions, represent the $L_{\rm [C\,II]}$ -SFR relations and 1σ uncertainties for low-redshift galaxies (De Looze et al. 2014) and highredshift ALPINE galaxies ($z \sim 4-6$; Schaerer et al. 2020), respectively. The dashed orange line shows the relation for low-z composite galaxies and AGNs (De Looze et al. 2014). Previous observations of z > 6 galaxies are plotted as gray points or downward triangles, which represent 3σ upper limits (Harikane et al. 2020; Schouws et al. 2023; Fudamoto et al. 2024).

nent with $T_{\rm dust}$ = 40 K results in a lower SFR and brings ID9094 into even better agreement with the $L_{\rm [CII]}$ -SFR_{tot} relation.

For ID2756, its 3σ upper limits of the [C II] emission line luminosity is $\log(L_{\rm [C\,II]}/L_{\odot}) < 8.58$ (Sect. 3.1). Due to the shallower NOEMA observations, we cannot constrain SFR well with and without AGN components in the SED (as described in Sect. 3.2.1). Therefore, we are unable to place meaningful constraints on its position in the $L_{\rm [C\,II]}$ versus SFR_{tot} plane without deeper FIR observations.

3.2.3. Caveats

Our analysis provides constraints on the dust and gas properties of the two LRDs. For ID9094, the two independent analyses of $T_{\rm dust}$ and the $L_{\rm [CII]}$ –SFR $_{\rm tot}$ relation show notable deviations from typical DSFGs. These distinct properties suggest that additional radiation mechanisms, such as AGN activity, contribute to the total emission of this source.

However, we emphasize that the interpretations above rely on several assumptions and modeling choices. First, the lower limit on the dust temperature ($T_{\rm dust} \gtrsim 110\,{\rm K}$) inferred in the absence of an AGN component is derived under fixed SED parameters and based solely on upper limits in the FIR. The resulting value is therefore model-dependent and sensitive to assumptions such as the dust emissivity index and IR SED shape. Moreover, high dust temperatures alone are not sufficient to

definitively rule out a purely star-forming scenario. Although rare, some high-redshift galaxies without clear AGN signatures have exhibited elevated $T_{\rm dust}$ values. For instance, MACS0416-Y1 at z=8.31 has a $T_{\rm dust}\sim80-116\,\rm K$, potentially due to intense starburst activity in a merger-driven system without requiring AGN heating (e.g., Bakx et al. 2020; Fudamoto et al. 2023; Harshan et al. 2024; Sommovigo et al. 2022). Similarly, A2744-YD4 at z=8.38 shows $T_{\rm dust}\sim55-107\,\rm K$, which may be due to its low metallicity and high SFR (e.g., Laporte et al. 2017; Behrens et al. 2018; Sommovigo et al. 2022). These cases are considered extreme and likely involve compact starbursts and low dust content.

Second, the observed offset from the empirical $L_{\rm [CII]}$ –SFR relation is not conclusive evidence against the DSFG scenario. The total SFR estimates are subject to uncertainties, especially when derived from upper limits, and the comparison sample includes a mix of UV-bright and IR-selected galaxies with diverse properties. Moreover, the $L_{\rm [CII]}$ –SFR relation exhibits large scatter at high redshift, and many star-forming galaxies also lie below the relation.

In summary, while our results are more consistent with an AGN-driven scenario, the possibility of dusty star formation remains viable given the current data limitations. Future deeper FIR observations will be essential to distinguish between these scenarios.

4. Impact of LRDs on our understanding of DSFGs

As introduced earlier, the two LRDs analyzed in this study were initially selected as the most massive DSFG candidates at $z_{\text{spec}} > 7$ in the JWST FRESCO GOODS-North field (Xiao et al. 2024). This misclassification arose due to the overlap in selection criteria between LRDs and DSFGs, particularly for optical/NIR-selected DSFGs often referred to as optically dark or faint galaxies (also known as HST-dark galaxies or H dropouts; e.g., Wang et al. 2019; Alcalde Pampliega et al. 2019; Williams et al. 2019; Xiao et al. 2023; Gómez-Guijarro et al. 2022; Barrufet et al. 2023; Akins et al. 2023; Pérez-González et al. 2023; Barro et al. 2024). Typically, DSFGs identified in optical/NIR observations are selected based on their red colors and rest-frame optical faintness, as these characteristics correlate strongly with dust attenuation at similar redshifts (e.g., Wang et al. 2016, Xiao et al. 2023). Additionally, the optical faintness criterion helps exclude quiescent galaxies, which are generally optically bright. For instance, DSFGs are often selected using a color cut such as F150W-F444W > 1.5 mag and $F150W \gtrsim 26.5$ mag, though specific band choices and thresholds slightly vary across studies.

With the increasing availability of JWST spectroscopic data, emission-line detections have become an additional criterion for DSFG selection (e.g., Xiao et al. 2024; Barrufet et al. 2025). In Xiao et al. (2024), DSFG candidates were selected based on their red color (F182M – F444W > 1.5 mag) and strong emission lines (e.g., $\text{H}\alpha$ + [NII] + [SII] lines or [O III] $\lambda\lambda$ 4960, 5008 + $\text{H}\beta$ lines). This approach also allows for the identification of LRDs among DSFG candidates. Among 26 DSFG candidates with detected $\text{H}\alpha$ lines (>8 σ), seven (~27%) were found to be LRDs, exhibiting broad $\text{H}\alpha$ emission lines (see Matthee et al. 2024).

Over the past decade, numerous studies have shown that optically dark or faint galaxies tend to have high stellar masses and may dominate the massive end of the stellar mass function (e.g., Wang et al. 2019; Gottumukkala et al. 2024). Additionally, these galaxies could contribute significantly to the cosmic star formation rate density (SFRD) at z > 3, with

contributions of ~10–50% of the SFRD, depending on the methodology (see a direct comparison in Xiao et al. 2023; e.g., Wang et al. 2019; Williams et al. 2019; Gruppioni et al. 2020; Fudamoto et al. 2021; Talia et al. 2021; Enia et al. 2022; Shu et al. 2022; Barrufet et al. 2023; van der Vlugt et al. 2023). Notably, while different studies employ varying selection criteria for optically dark or faint galaxies, some adopt stricter definitions, requiring FIR/submillimeter (e.g., ALMA/NOEMA) detection and/or non-detection in HST. Given that LRDs have been found to 1) lack FIR/submillimeter detections (see also, e.g., Labbe et al. 2025; Williams et al. 2024), 2) exhibit a characteristic "v-shaped" SED (indicating detection by the HST), they are unlikely to contaminate the DSFG sample in these studies.

On a global scale, these findings suggest that LRD contamination should not significantly bias SFRD estimates derived from DSFG samples. However, their presence may still impact number density estimates of optical/NIR-selected DSFG populations, particularly when broader color-based selection criteria are used. A study with a larger sample is needed to determine how much the number density estimates of DSFG populations are affected by LRDs. These results highlight the importance of refining DSFG selection methods to minimize potential contamination from AGN-dominated LRDs, ensuring a more accurate census of dust-obscured star formation in the early Universe.

5. Conclusions

In this paper we present a study of two LRDs, ID9094 and ID2756, recently discovered in the JWST FRESCO GOODS-North field at $z_{\rm spec} > 7$. These sources, characterized by their red, point-source-like morphology and "v-shaped" SEDs, exhibit broad H β emission lines, which suggests they are AGN candidates. Using NOEMA observations targeting the [C II] 158 μ m emission line and the 1.3 mm dust continuum, we obtain the following key results:

- Both sources are undetected in the [C II] and 1.3 mm continuum. For ID9094, the 3σ upper limits on the [C II] and 1.3 mm continuum are far below expectations for a typical DSFG, which would have been detected at >16 σ in [C II] and >3 σ at 1.3 mm. For ID2756, the expected >5 σ detection in [C II] is also not observed.
- Benefiting from the 1.3 mm upper limits, we performed UV-to-FIR SED fitting, which provided the first hint of an AGN scenario in this work. For ID9094, without including an AGN component, the SED fit requires an abnormally high dust temperature ($T_{\rm dust} \gtrsim 110~{\rm K}$), significantly above typical values for star-forming galaxies at similar redshifts. Including an AGN component yields more physically plausible fits with dust temperatures that can be as low as ~40–60 K. However, we caution that the minimum $T_{\rm dust}$ is strongly linked to the AGN fraction contributing to the total IR luminosity higher AGN fractions allow for lower $T_{\rm dust}$ values. For ID2756, the SED fitting remains inconclusive due to shallower NOEMA data, highlighting the need for deeper FIR/submillimeter observations.
- Using the [C II] upper limits, we compared the sources to the empirical L_[C II]-SFR_{tot} relation; this provides a second hint of an AGN scenario. For ID9094, its [C II] upper limit places it significantly below the scaling relation when no AGN component is included, but consistent with the relation when an AGN contribution is considered. The case for ID2756 remains uncertain due to the limited NOEMA depth.

These results highlight the value of FIR constraints, when combined with shorter-wavelength photometry, in revealing the

nature of LRDs – particularly when spectroscopic data are unavailable. For ID9094, the results are consistent with a scenario in which AGN activity contributes to the observed emission, though a dusty star-forming origin cannot be fully excluded given the data limitations. We also note that both sources are detected in JWST/MIRI F1000W, and ID9094 is additionally detected in F2100W, evidence in support of the presence of hot dust emission (G. Leung, priv. comm.). For ID2756, the current observations are insufficient to determine the dominant energy source. We emphasize that our interpretation remains modeldependent and subject to assumptions in the SED fitting and limited FIR constraints. More broadly, this study demonstrates the critical role of FIR data in clarifying the nature of LRDs. Future deeper FIR/submillimeter campaigns, along with highresolution spectroscopy and multiwavelength data, will be essential to fully understanding the physical properties of these LRDs and their role in early galaxy evolution.

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Appendix A: Best-fit SED of ID2756 with $T_{\rm dust} = 40$ K.

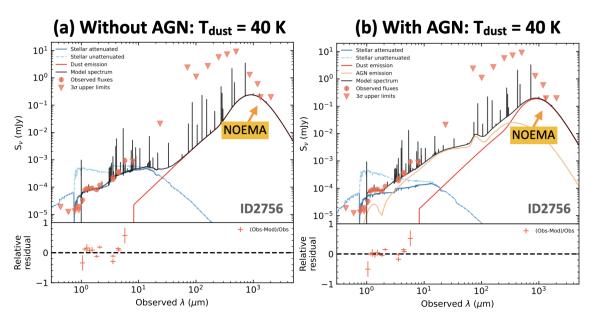


Fig. A.1. Similar to Fig. 4 but for ID2756. The 3σ upper limit at 1.3 mm is not sufficiently deep to rule out the low $T_{\rm dust}$ solution, nor to rule out the scenario without an AGN. Deeper FIR/submillimeter observations, at least similar to those for ID9094 (rms level $\sim 30.4~\mu \rm Jy~beam^{-1}$ at 1.3 mm), are needed.